

# CARS Measurements of Vibrational Distribution Function in Nanosecond Pulsed Plasmas

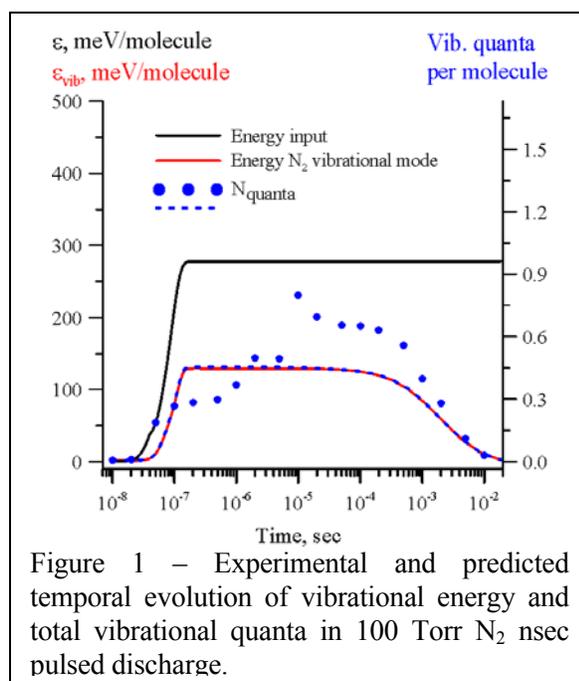
Walter R Lempert and Igor V. Adamovich

Departments of Mechanical & Aerospace Engineering and Chemistry, The Ohio State University, Columbus, OH, lempert.1@osu.edu

Picosecond Coherent Anti-Stokes Raman Scattering (CARS) has been used for studies of vibrational energy dynamics in a diffuse single filament pin-to-pin discharge in pure N<sub>2</sub> and air. The discharge produces a filament with approximate dimensions 2-3 mm (diameter) x 10 mm (height) using a pair of 7.5 mm bare copper spherical electrodes.

Results for N<sub>2</sub> at P = 100 Torr are summarized in Fig. 1, which shows experimental and Master equation modeling predictions for total vibrational quanta/molecule in levels v = 0 – 8 for which population was observed, as well as discharge total specific energy loading, and modelling predictions for energy stored in vibration. It can be seen that experimentally approximately 30% of coupled pulse energy is loaded into the N<sub>2</sub> vibrational mode which is in qualitative agreement

with the modelling prediction of approximately 45%. However it can also be seen that the total vibrational quanta/molecule rises rapidly after the termination of the discharge pulse, by a factor of between ~two and three for time delays after discharge initiation in the range ~10–200 microseconds. This suggests that additional energy, up to a maximum of ~70% of the coupled discharge pulse energy, is loaded into the vibrational mode after the discharge pulse. This trend is clearly at variance with the kinetic model, which predicts the number of quanta to remain nearly steady for up to t~200 μsec after the pulse, since V-V exchange among nitrogen molecules conserves the total number of quanta in the N<sub>2</sub> vibrational mode. At longer time delays slow relaxation is observed due to a combination of V-T relaxation, mass diffusion, and convection.



The results summarized in Fig.1, as well as rotational/translational temperature and more detailed measurements of the temporal evolution of the nitrogen vibrational distribution function, suggests that a significant fraction of energy initially stored in excited electronic species and dissociation products by electron impact is subsequently transferred to the vibrational mode of the ground N<sub>2</sub>(X<sup>1</sup>Σ, v), rather than to heat, during collisional processes such as the energy pooling process, N<sub>2</sub>(A) + N<sub>2</sub>(A) → N<sub>2</sub>(B, C) + N<sub>2</sub>(X, v). Similar results are obtained in air at P = 100 Torr and from spontaneous Raman scattering measurements in which detectable population was observed in vibrational levels as high as v = 12.

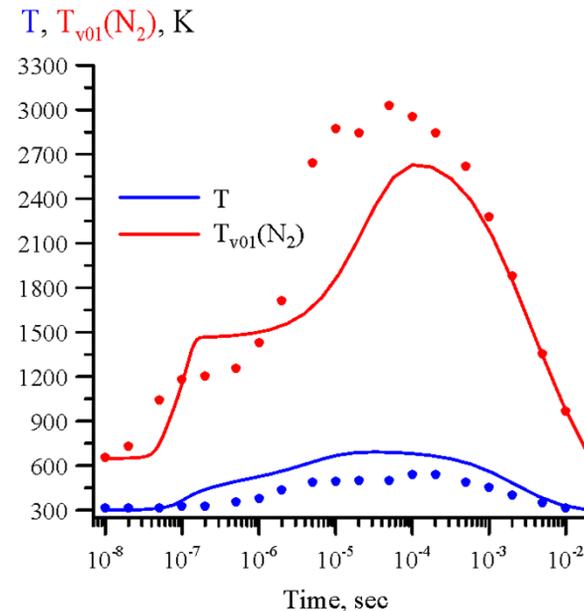
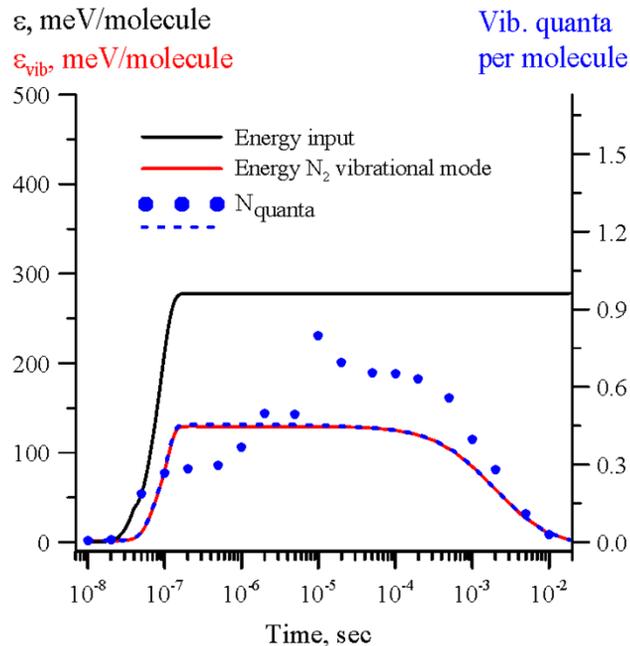
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**Highlight**



# ENERGY LOADING AND DECAY OF N<sub>2</sub> VIBRATION IN PULSED N<sub>2</sub> AND AIR PLASMAS: NEW PHYSICS?

- CARS (Coherent Anti-Stokes Raman Scattering) laser diagnostics demonstrated that ~40% of coupled energy from ~100 ns discharge pulse is initially loaded into N<sub>2</sub> vibrations (v=0-8 detected) by electron impact.
- N<sub>2</sub> vibrational energy increases to ~80% of coupled energy for time scales ~10-200 μsec after discharge pulse (symbols). Model (solid lines) predicts vib energy to be constant, under-predicts 1<sup>st</sup> level vib temperature, and over-predicts rotational temperature.
- Mechanism for N<sub>2</sub>(X,v) population after the pulse is unknown – quenching of N<sub>2</sub>(A,B,C,a) quenching or N atom recombination are possibilities.



- 1<sup>st</sup> Level Vibrational and Rotational Temperature

- Vibrational Energy of N<sub>2</sub> (v=0-8)

# Particle Charging in Argon-Hydrogen Plasmas

Meenakshi Mamunuru and Uwe R. Kortshagen

Department of Mechanical Engineering, University of Minnesota, Minneapolis, MN  
(Kortshagen@umn.edu)

Capacitively coupled plasma (CCP) RF discharges are used for gas phase synthesis of silicon nanocrystals. Typically,  $\text{SiH}_4$  diluted in  $\approx 95\%$  argon is flowed through the discharge chamber at a pressure of a few Torr [1]. The plasma dissociates  $\text{SiH}_4$ , creating radicals that react and form uniformly sized, crystallized silicon nanoparticles which are collected downstream. Interaction of the plasma with the nanoparticles plays a role in determining their surface and bulk properties. The particles acquire a steady state negative charge and are impacted by energetic ions in the plasma. Processes such as rearrangement of atoms in the particle to form a crystal lattice and termination of the atoms on the surface depend on the impaction energies of the ions.

Trace amounts of  $\text{H}_2$  ( $\approx 2\%$ ) are produced in the discharge due to dissociation of  $\text{SiH}_4$  and subsequent gas phase reactions. The  $\text{H}_2$  is ionized and the plasma may consist of comparable fractions of heavy Ar ions and light hydrogen ions [2]. Plasma-particle interactions in the presence of comparable fractions of argon and hydrogen ions were investigated.

An analytical model was developed to calculate particle charging in plasmas and the results of the model agree with Monte Carlo simulations. The steady state particle potential, ion flux and average ion impact energy for a 500 nm particle are shown in Fig. 1. The discharge consists of either only  $\text{Ar}^+$  ions or equal densities of  $\text{Ar}^+$  and  $\text{H}_3^+$  ions. In the pure argon discharge the particle potential reduces (becomes less negative) at intermediate pressures due to enhancement of the ion current, and increases (more negative) at higher pressures as the ion current is inhibited by collisions. In the presence of both  $\text{Ar}^+$  and  $\text{H}_3^+$  ions, the ion flux of each species to the nanoparticle peaks at different pressures due to the difference in their ion-gas collision characteristics. The  $\text{H}_3^+$  ion current peaks at a higher pressure compared to the  $\text{Ar}^+$  current. The presence of  $\text{H}_3^+$  ions causes a reduction in the particle potential (less negative) over a wide pressure range, and so a reduction in the average ion impact energy.

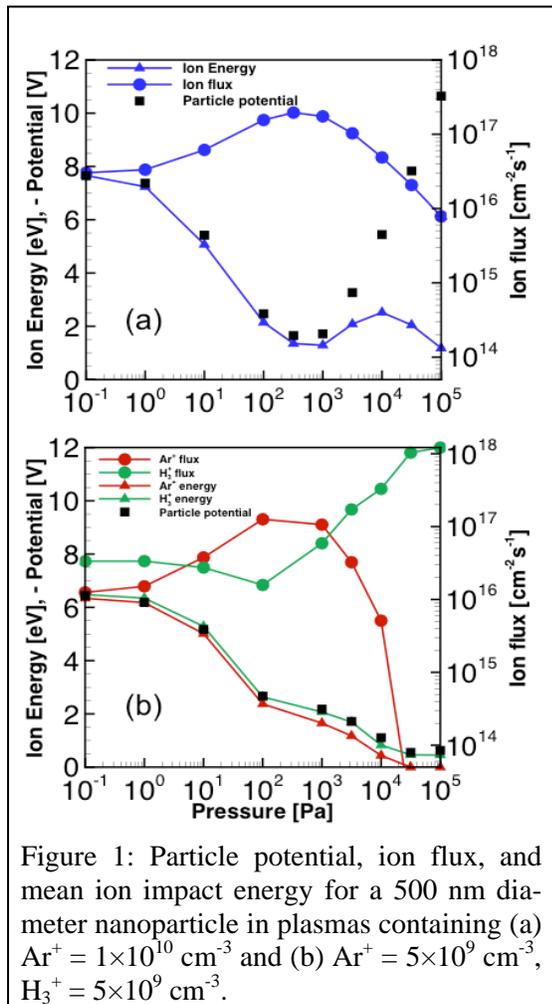


Figure 1: Particle potential, ion flux, and mean ion impact energy for a 500 nm diameter nanoparticle in plasmas containing (a)  $\text{Ar}^+ = 1 \times 10^{10} \text{ cm}^{-3}$  and (b)  $\text{Ar}^+ = 5 \times 10^9 \text{ cm}^{-3}$ ,  $\text{H}_3^+ = 5 \times 10^9 \text{ cm}^{-3}$ .

## References

- [1] L. Mangolini, Nano Letters **5**, 4 (2005).
- [2] A. Bogaerts, and R. Gijbels, Spectrochimica Acta Part B: Atomic Spectroscopy **57**, 6 (2002).

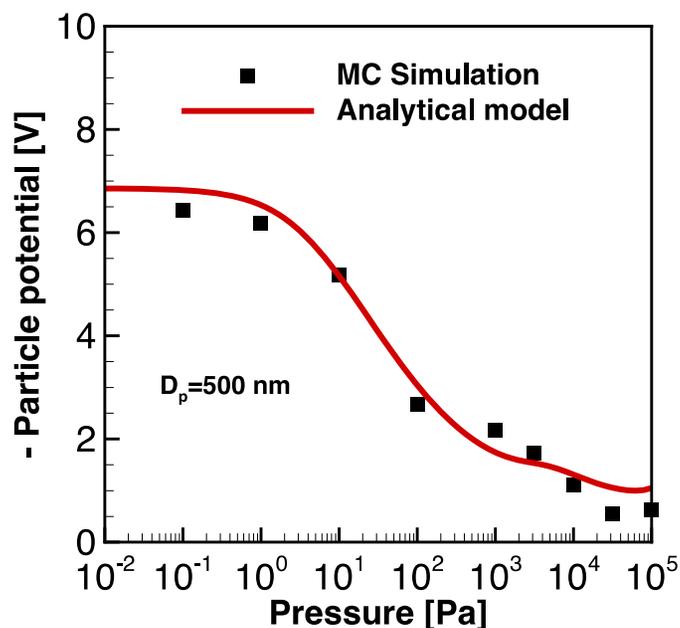
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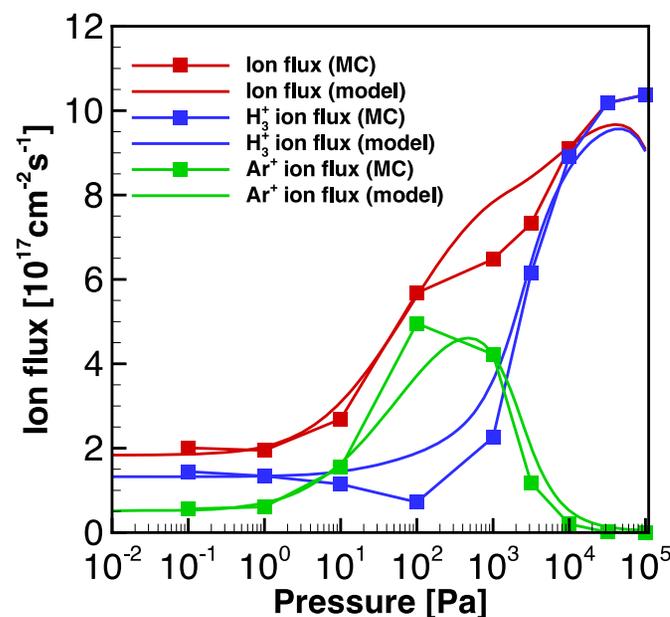


# PARTICLE CHARGING IN Ar-H<sub>2</sub> PLASMAS

- An analytical model calculates steady state nanoparticle charge in a plasma containing multiple positive ions and compares well with Monte Carlo simulations.
- In Ar-H<sub>2</sub> discharges containing equal number of Ar<sup>+</sup> and H<sub>3</sub><sup>+</sup> ( $5 \times 10^9 \text{ cm}^{-3}$ ) ions, flux of each species to the particle depends on the individual ion-gas collision characteristics and so fractional current is pressure dependent.



- Steady state particle potential



- Total and individual ion fluxes to particle