

# Control of the Ion Energy Distribution on a Plasma Electrode

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# Outline

- **Introduction**
- **PIC-MCC simulation of tailored bias on boundary electrode**
- **Semi-analytic model**
- **Comparison with data**
- **Summary**



## Introduction / Motivation

- **Control of the energy of ions bombarding a substrate in contact with plasma is critical for plasma processing.**
- **The ion energy must be high enough to drive anisotropic etching, but not too high to induce substrate damage and/or loss of selectivity.**
- **As device dimensions continue to shrink, precise control of the ion energy *distribution* (not just the average ion energy) becomes increasingly important.**

## Goal and Approach

### Goal:

Develop methodologies to achieve “tailored” ion energy distributions<sup>1</sup> (IEDs).

### Approach:

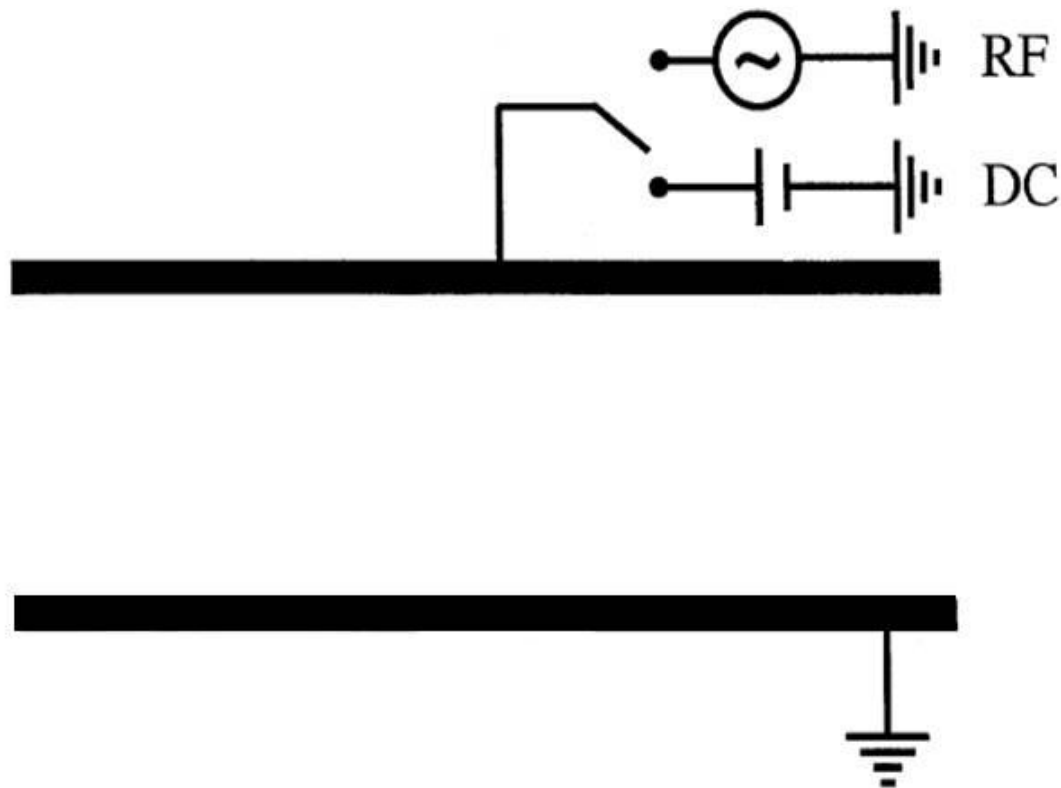
- Use combination of modeling/simulation and experiments.
- Particle-In-Cell simulation with Monte Carlo Collisions (PIC-MCC).<sup>2</sup>
- Semi-analytic model<sup>3</sup>.

<sup>1</sup>P. Diomede et al., *J. Appl. Phys.*, **109**, 083302 (2011); X.V. Quin, Y.-H. Ting and A.E. Wendt, *Plasma Sources Sci. Technol.*, **19**, 065014 (2010).

<sup>2</sup>S. K. Nam, D. J. Economou and V. M. Donnelly, *IEEE Trans. Plasma Sci.*, **35**, 1370 (2007); V. Vahedi, G. DiPeso, C. K. Birdsall, M. A. Lieberman, and T. D. Rognlien, *Plasma Sources Sci. Technol.*, **2**, 261 (1993); J. Verboncoeur, M. Alves, V. Vahedi, and C. K. Birdsall, *J. Comp. Phys.*, **104**, 321 (1993); P. Diomede et al. *PSST*, **14**, 459 (2005)

<sup>3</sup>P. Diomede et al., *PSST*, **20**, 045011 (2011).

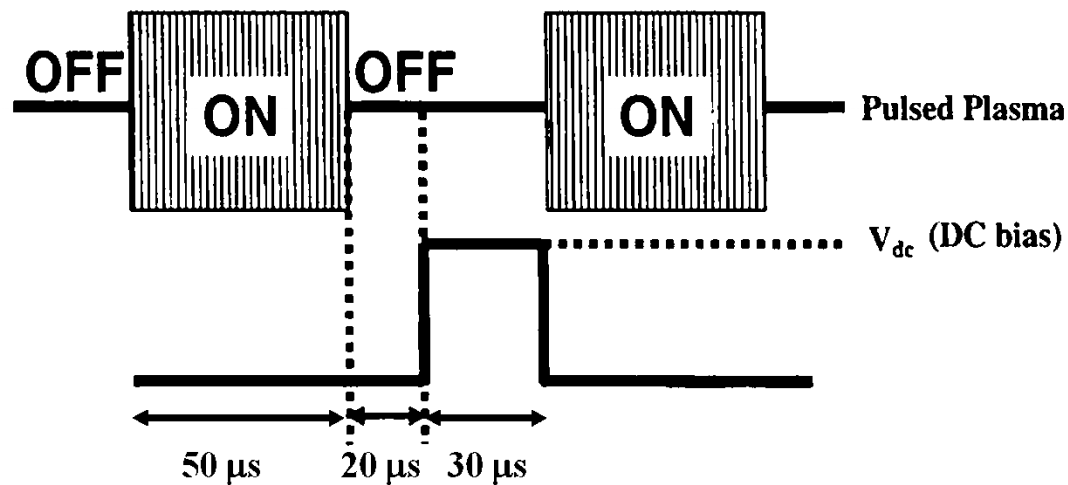
# PIC-MCC Simulation of Pulsed CCP Reactor with DC Bias in Afterglow



Argon plasma  
 $V_{\text{RF}} = 300 \text{ V}$   
 $\nu_{\text{RF}} = 13.56 \text{ MHz}$   
 $p = 10 \text{ mTorr}$   
 $d = 6 \text{ cm}$   
  
10 kHz pulse  
50% duty ratio

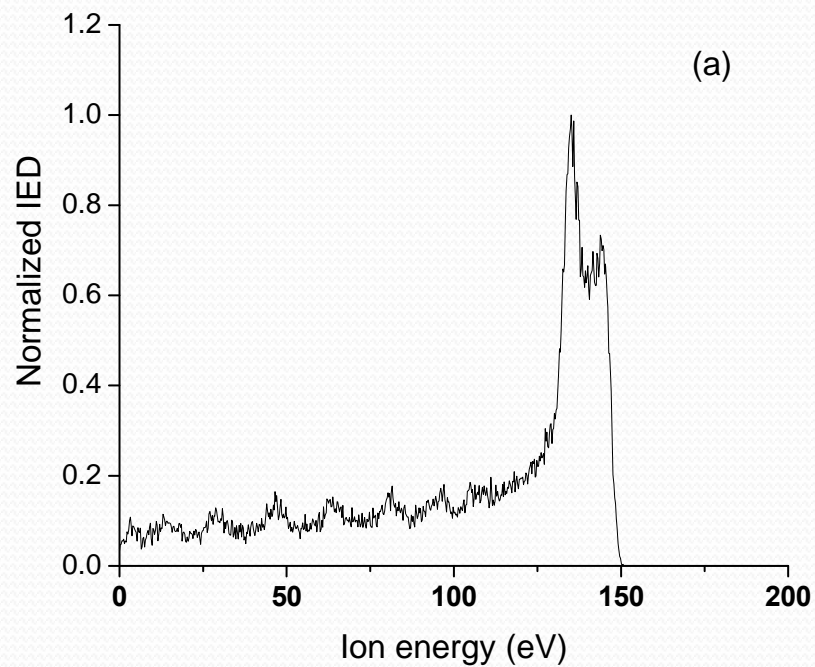
- Pulsed plasma is sustained in capacitively coupled plasma (CCP) reactor.
- DC bias is applied on the upper (boundary) electrode in the afterglow to modify the IED on the lower (substrate) electrode.

# Application of DC Bias in the Afterglow of a Pulsed Plasma

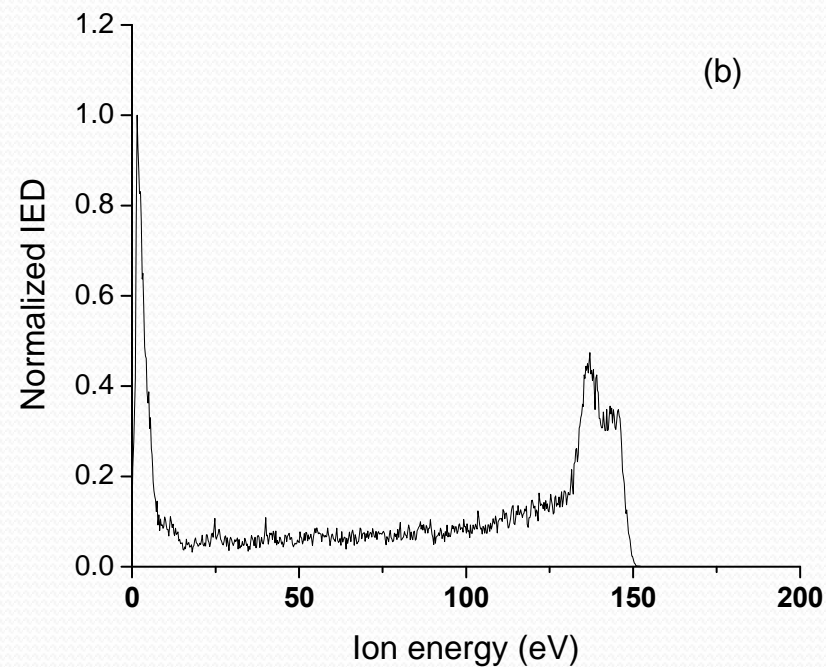


- After plasma power turn off (afterglow),  $T_e$  and  $V_p$  decay rapidly.
- Apply synchronously tailored positive bias voltage  $V_{dc}$  during specified time window in the afterglow.
- Bias raises plasma potential, modifying the IED on the wafer.

# IED without Bias

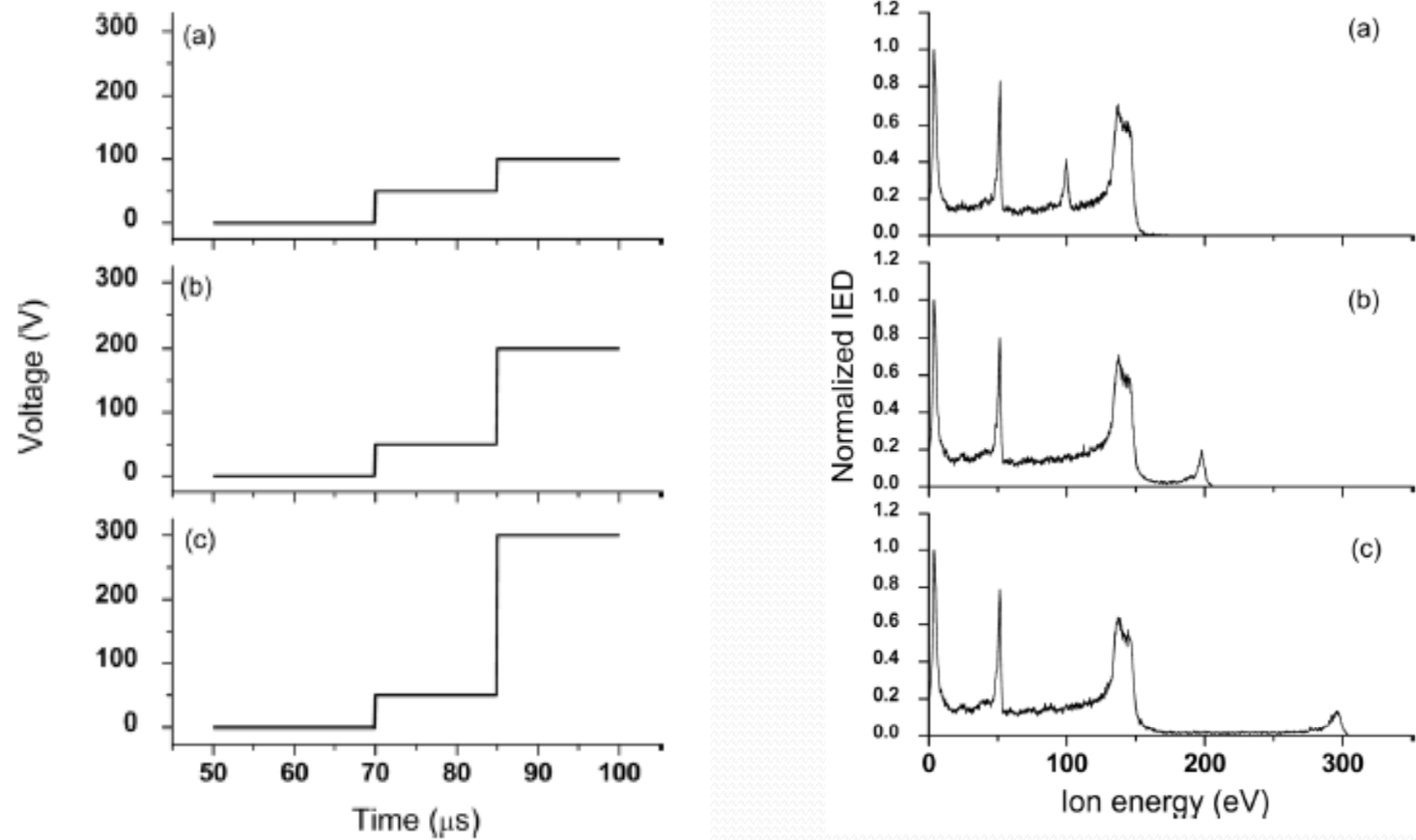


(a) IED for continuous wave (cw) plasma w/o bias



(b) IED for pulsed plasma (10 kHz, 50% duty cycle) w/o bias

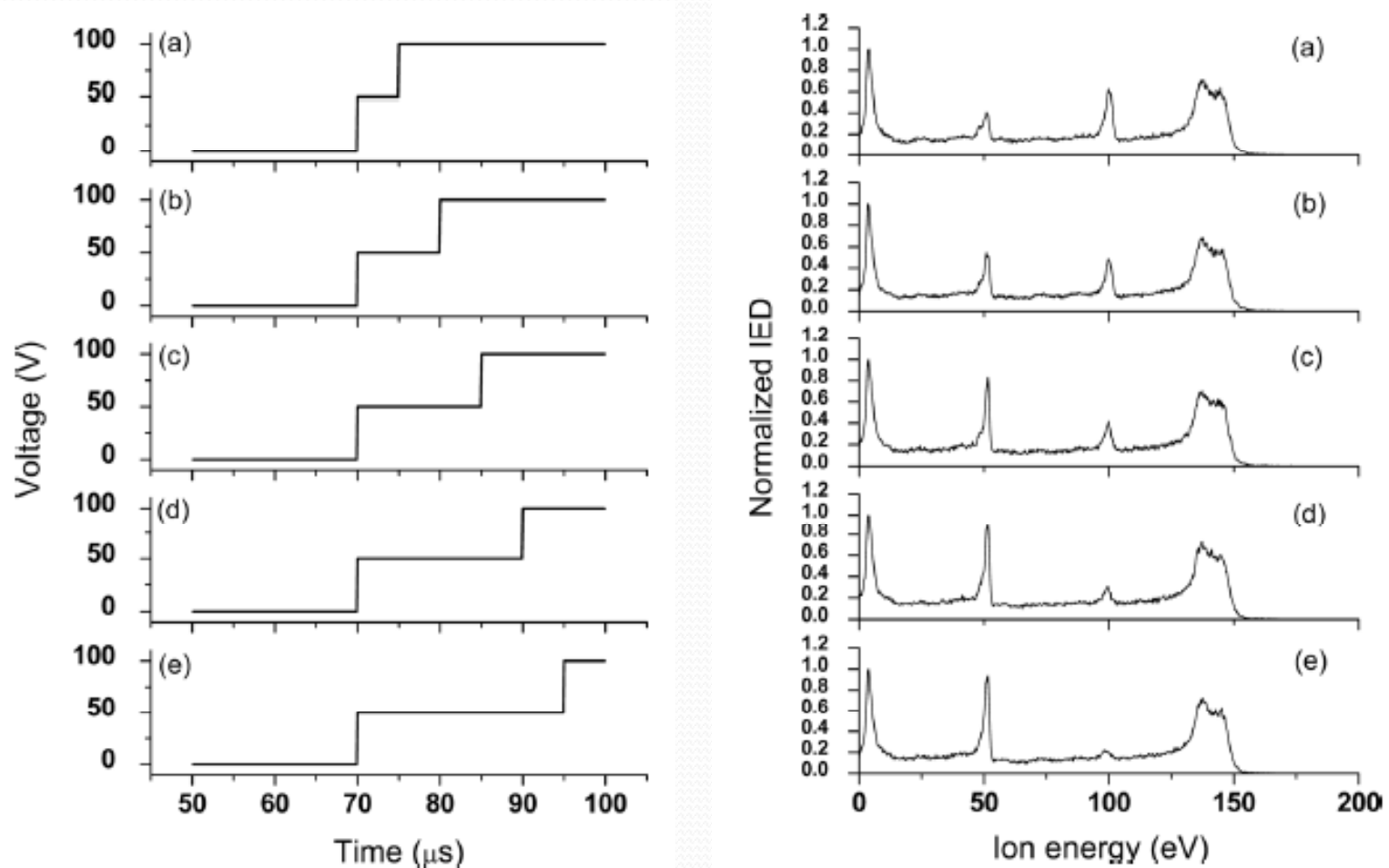
# IEDs with Staircase DC Bias Applied in Afterglow



- Afterglow starts at time  $t = 50 \mu\text{s}$ .
- Additional peaks appear in the IED.
- Peak location can be controlled by the value of the applied bias voltage.



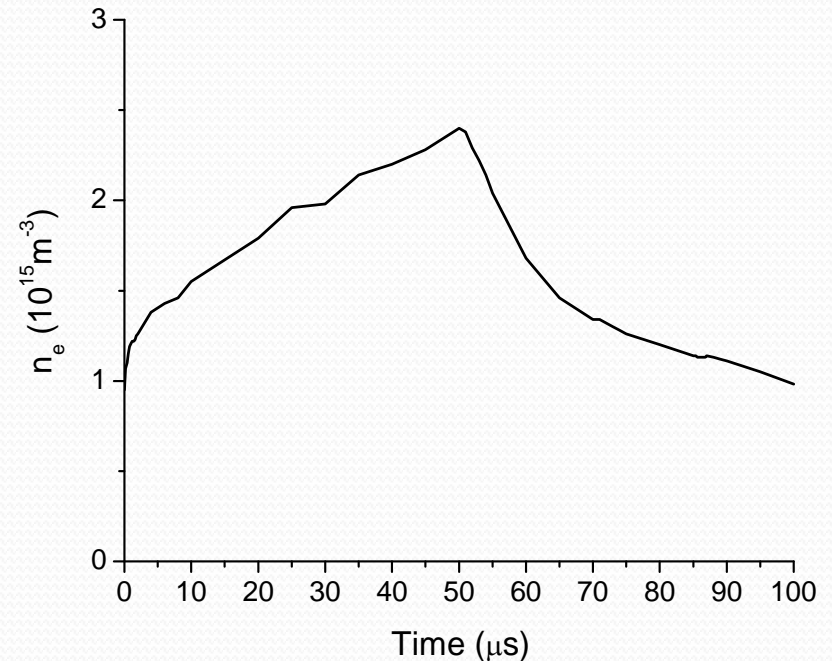
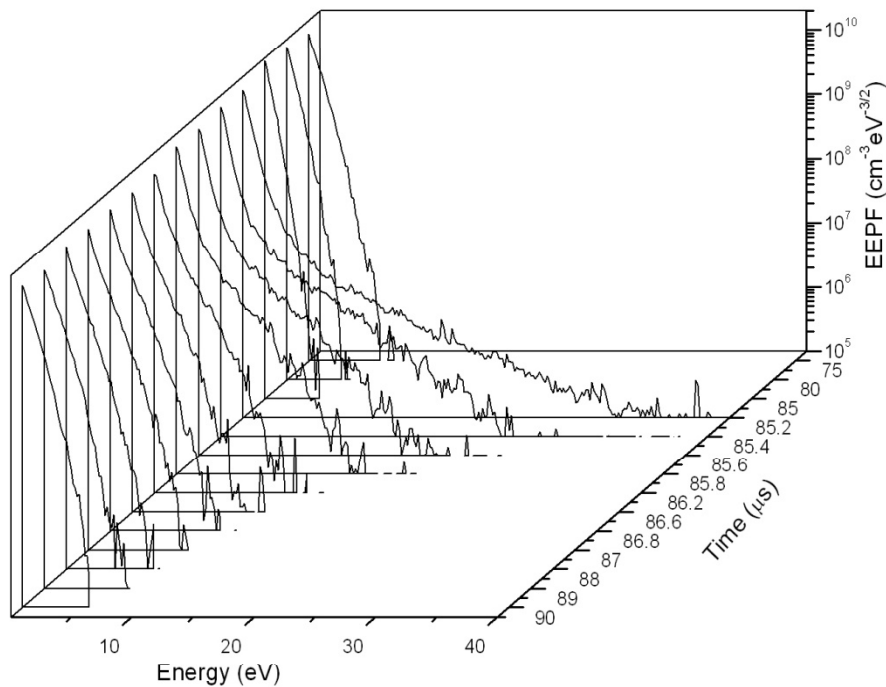
## IEDs with Staircase DC Bias Applied in Afterglow (2)



Peak strength can be controlled by the duration of the respective DC bias voltage.

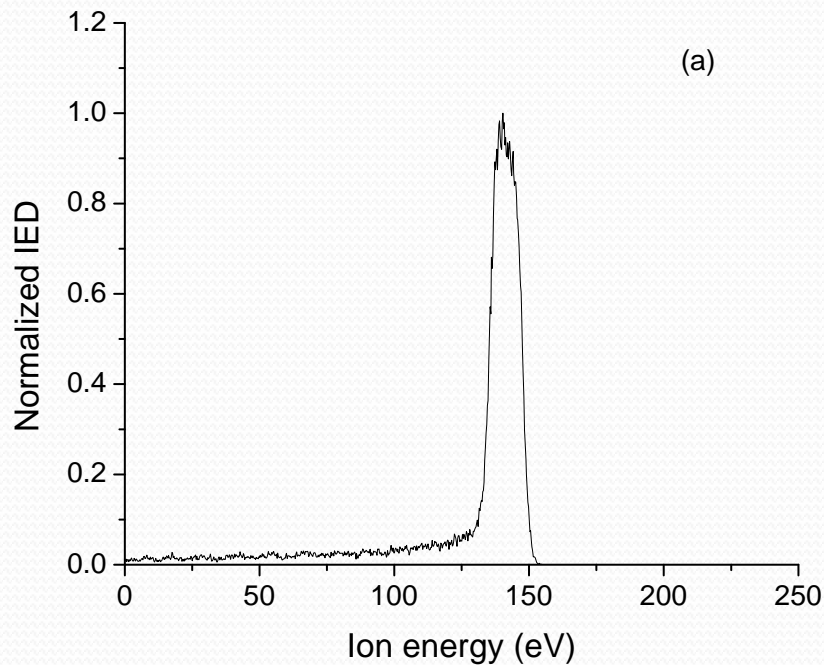
P. Diomede et al., *J. Appl. Phys.*, **109**, 083302 (2011)

# EEPF and Electron Density Evolution

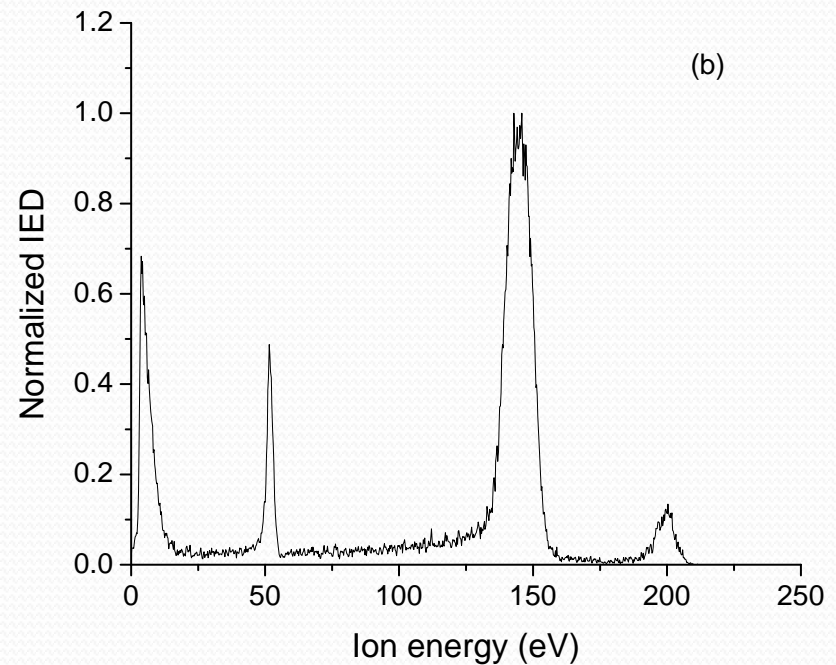


- In the afterglow of pulsed CCP, apply 50 V DC during  $t = 70-85 \mu\text{s}$ , followed by 300 V DC during  $t = 85-100 \mu\text{s}$ .
- EEPF is temporarily heated when the bias is applied but the electron density evolution is hardly affected.

# Plasma Excitation Frequency of 60 MHz



IED for cw plasma w/o bias



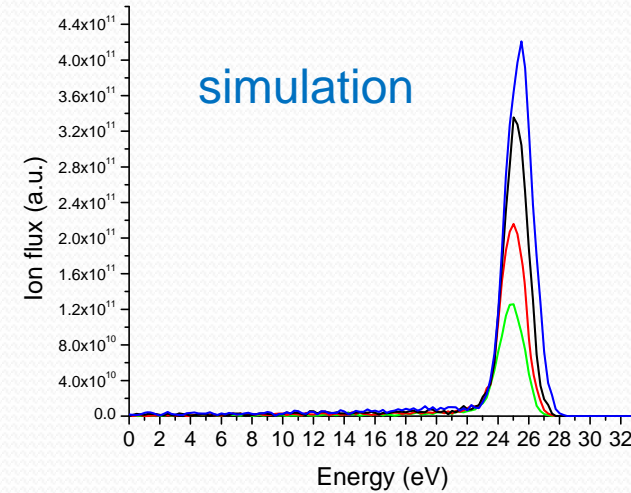
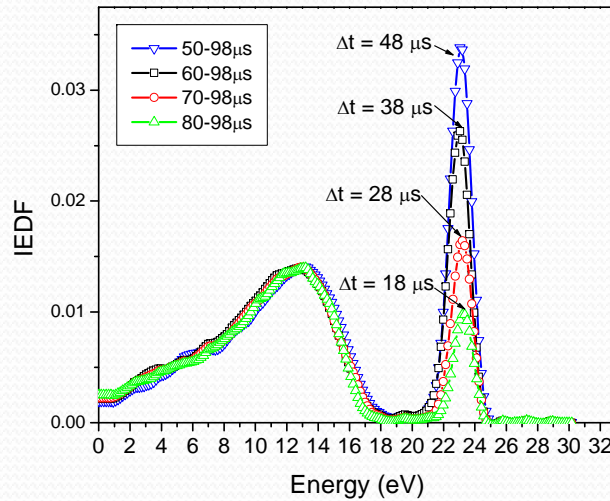
IED for pulsed plasma with staircase DC bias applied in afterglow (50 V during 70-85  $\mu$ s, 200 V during 85-100  $\mu$ s).

- Increasing frequency to 60 MHz boosts the electron density ( $n_e \sim f^2$ ), yields thinner sheath, reduces the collisionality of the sheath, and results in sharper IEDs.
- Ion energy peak location and fraction of ions under each peak can be controlled.

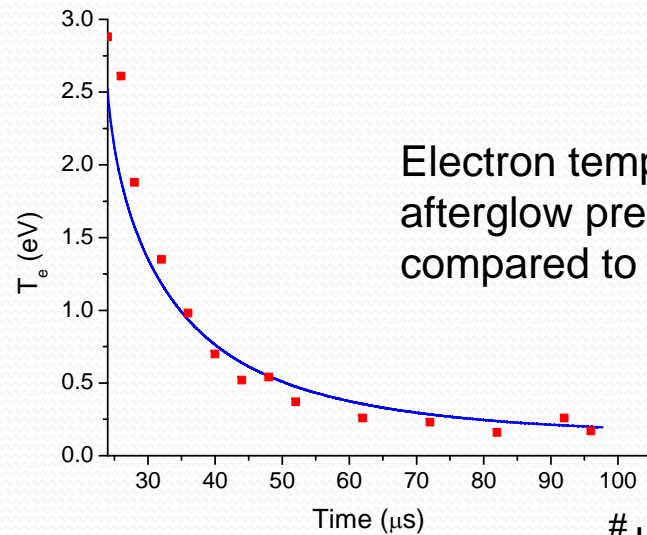
# Comparison of PIC Simulation with Experimental Data (1)

PS2-TuA9

PS+SS-WeM6



IEDs predicted by the PIC simulation of the afterglow (right) compared to data<sup>#</sup> (left). The low energy peak of the data is due to the active glow (not simulated by PIC).

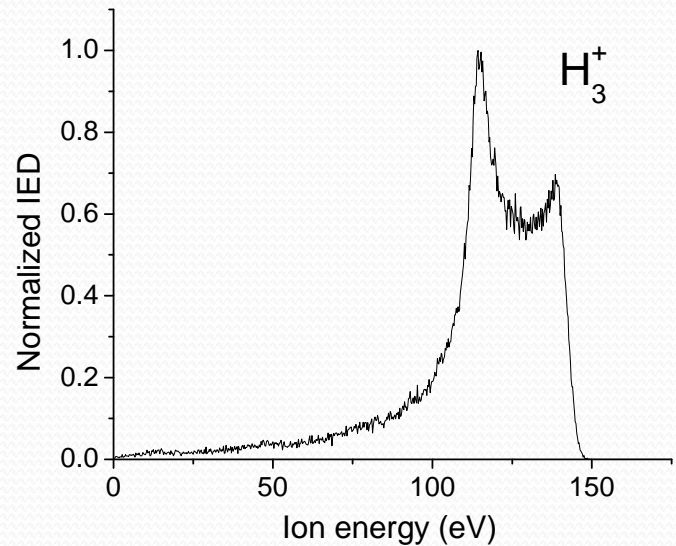
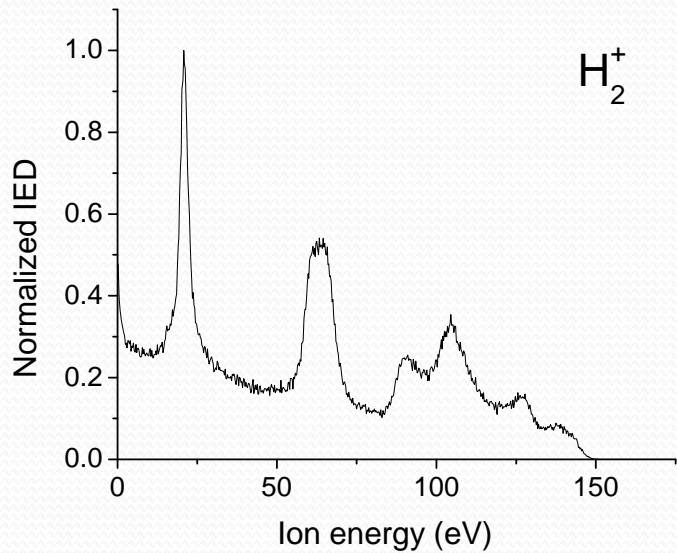


Electron temperature in the afterglow predicted by PIC (line), compared to data<sup>#</sup> (points).

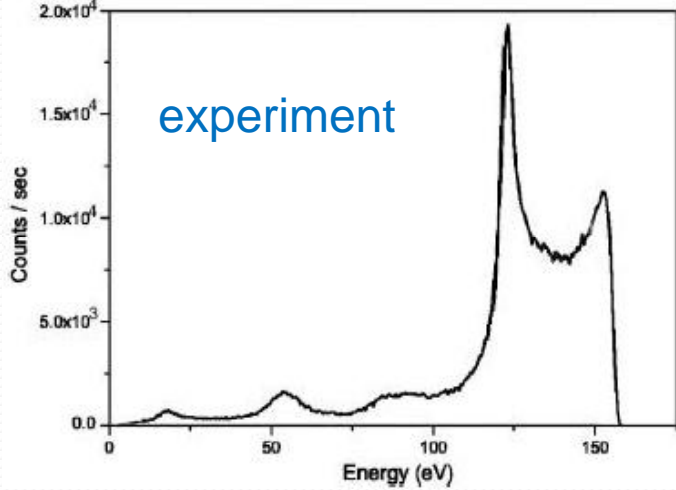
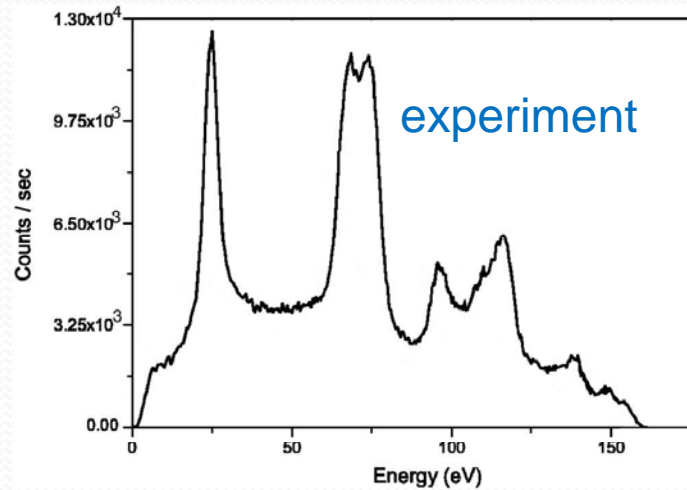
**Pulsed plasma, 10 KHz modulation, 20% duty, 14 mTorr, 120 W average power, 24.4 V DC bias applied in afterglow during time windows shown above.**

<sup>#</sup> H. Shin et al., *PSST*, **20** 055001 (2011).

# Comparison of PIC Simulation with Experimental Data (2)



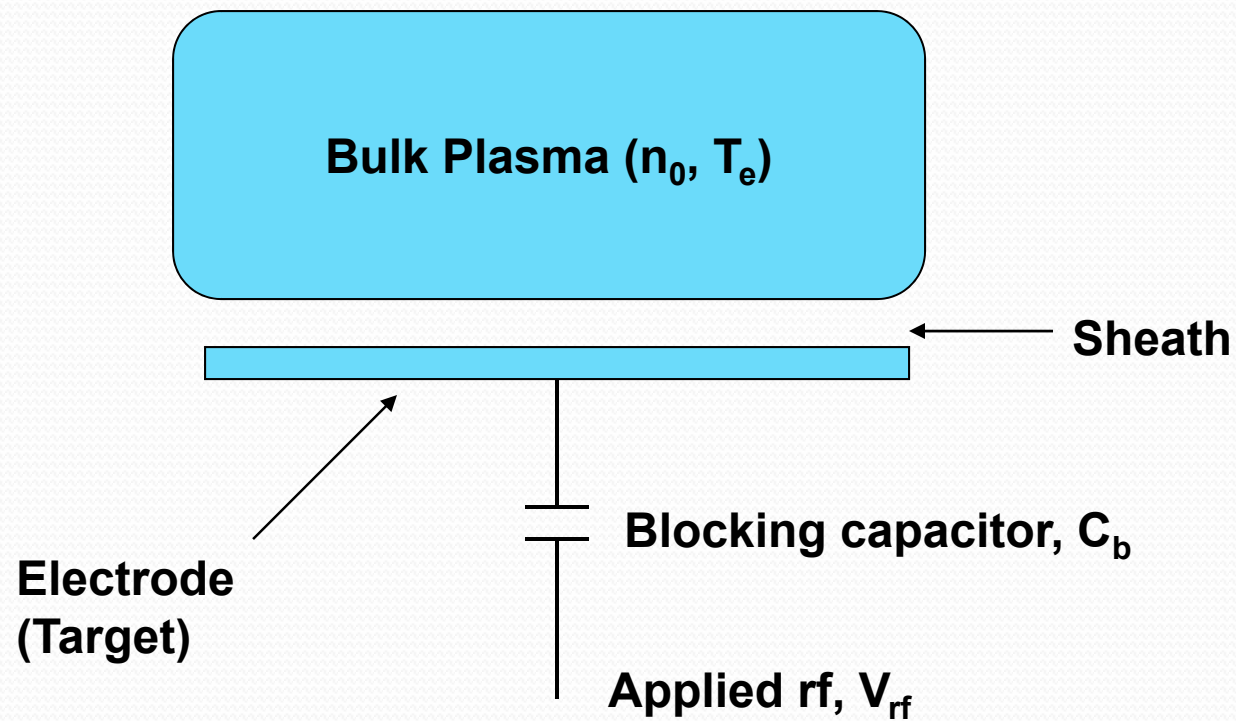
Hydrogen plasma  
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IEDs for cw plasma w/o bias

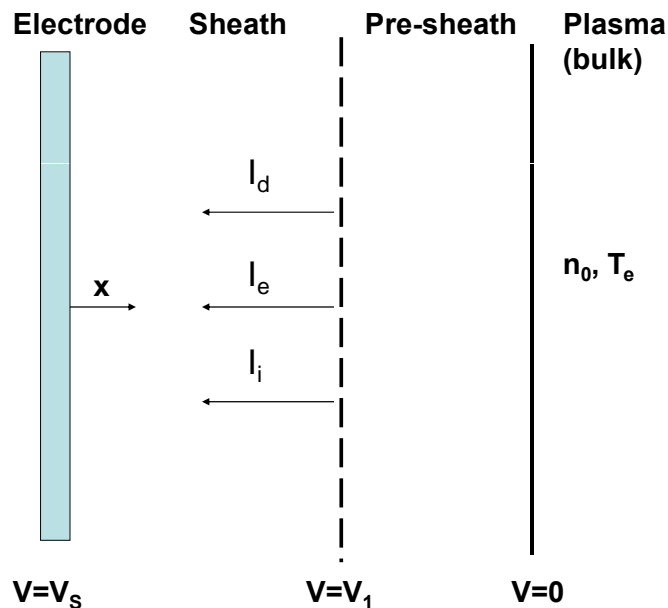
Experiments: D. O'Connell et al. *Phys. Plasmas* **14**, 103510 (2007)  
 Model: P. Diomede et al. *PSST*, **14**, 459 (2005)

## Electrode in Contact with Plasma



## Semi-analytic Model (1)

### Schematic of the sheath region



1. Electrode immersed in semi-infinite plasma of given electron (ion) density and electron temperature.
2. Electron, ion and displacement currents flow through the sheath.
3. Non-linear sheath capacitance  $C_s$  is calculated from the electric field at the electrode,  $E$ .

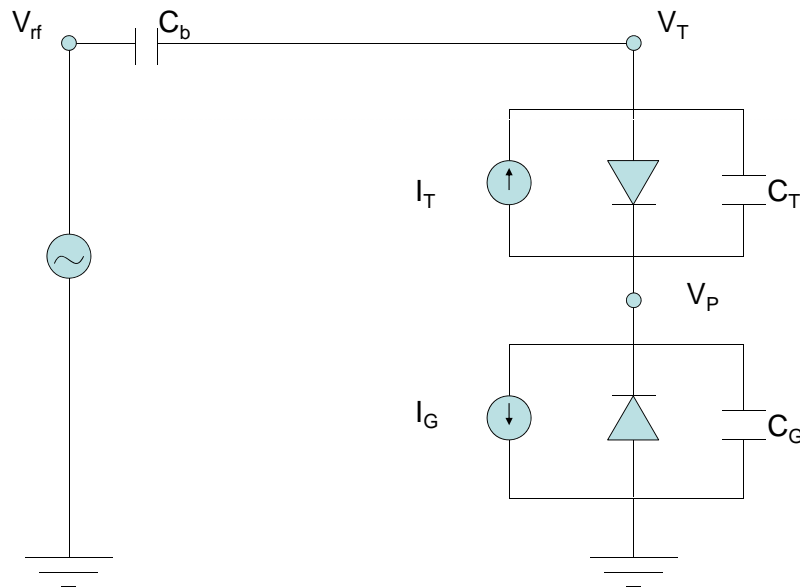
$$C_s = -\epsilon_0 A \frac{\partial E}{\partial V_s}$$

$$E = -\sqrt{\frac{2n_1 k T_e}{\epsilon_0}} \left[ \exp\left(\frac{e(V_s - V_1)}{k T_e}\right) + \frac{V_s}{V_1} - 2 \right]^{1/2}$$

- A. Metze et al., *J. Appl. Phys.*, **60**, 3081 (1986).  
 P. Miller and M. Riley, *J. Appl. Phys.*, **82**, 3689 (1997).  
 T. Panagopoulos and D. Economou, *JAP*, **85**, 3435 (1999).

## Semi-analytic Model (2)

Equivalent circuit model, A. Metzger et al., *J. Appl. Phys.*, **60**, 3081 (1986).



Subscripts T and G refer to “target” and “ground” electrodes, respectively.

$$C_b \frac{d}{dt}(V_{rf} - V_T) + C_T \frac{d}{dt}(V_P - V_T) + I_T = 0$$

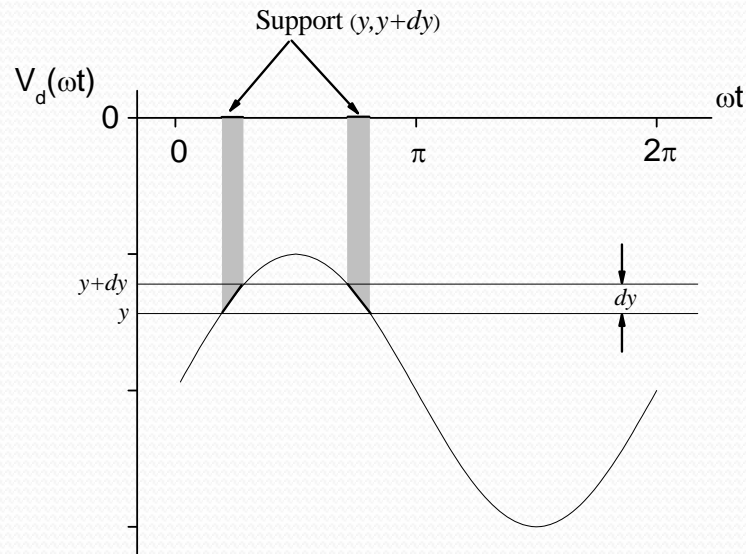
$$C_T \frac{d}{dt}(V_P - V_T) + C_G \frac{d}{dt}V_P + I_T + I_G = 0$$

$$\frac{dV_d}{dt} = -\frac{V_d - (V_T - V_P)}{\tau_i} \quad \text{Ions respond to a “damped” potential } V_d$$

Voltage  $V_{rf}$  is applied through blocking capacitor,  $C_b$ .  
**Given  $n_0$ ,  $T_e$ ,  $V_{rf}$  and  $C_b$ , calculate  $V_T$ ,  $V_P$ , and  $V_d$ .**



# Ion Energy Distribution



$$f(y) = \frac{1}{2\pi} \sum_{\substack{\# \text{ of points in } 0 < \omega t < 2\pi \\ \text{such that} \\ V_d(\omega t) = y}} \left| \frac{dV_d}{d(\omega t)} \right|$$

$$f(y) = IED$$

$$V_d = V_d(\omega t) = \text{"damped" sheath voltage}$$

Sample damped sheath potential waveform

$$y = V_d(\omega t) \Rightarrow \omega t = V_d^{-1}(y)$$

$$f(y) = \frac{1}{2\pi} \sum_{\substack{\# \text{ of points in } 0 < \omega t < 2\pi \\ \text{such that } V_d(\omega t) = y}} \left| \frac{dV_d^{-1}(y)}{dy} \right|$$

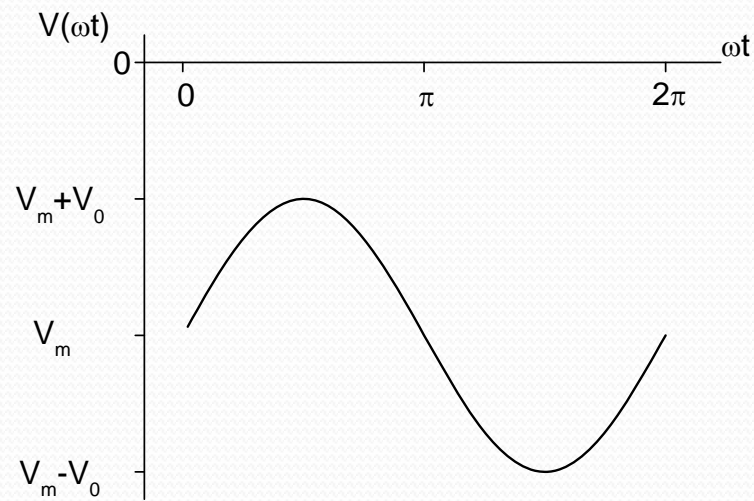
P. Diomede et al., *Plasma Sources Sci. Technol.*, **20**, 045011 (2011).

E. Kawamura et al., *Plasma Sources Sci. Technol.*, **8**, R45 (1999).

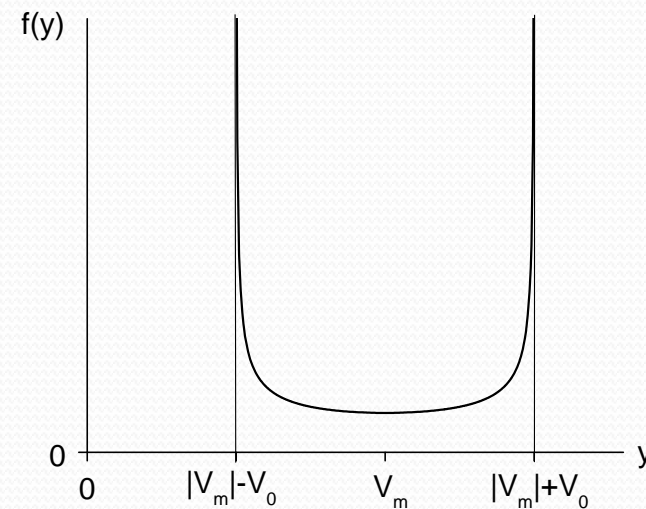
# IED for a sinusoidal sheath voltage

(Forward problem)

Damped sheath potential



Resulting IED



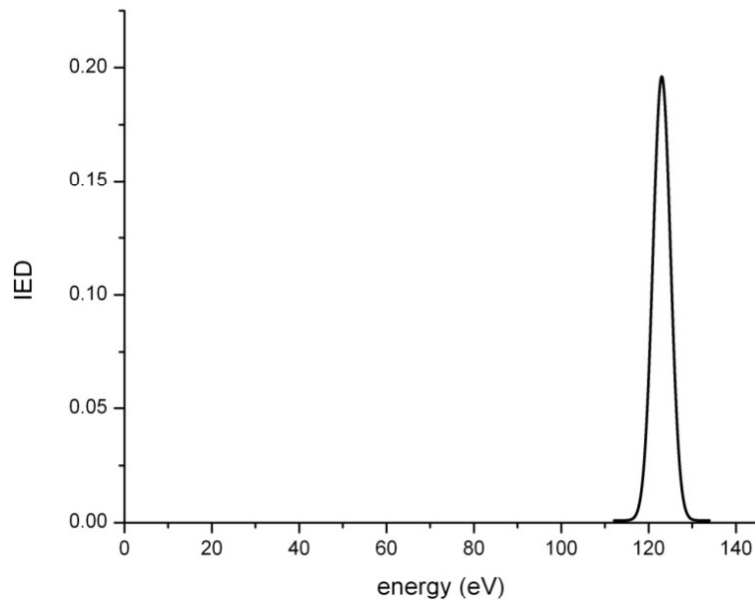
$$y = V_d(\omega t) = V_m + V_0 \sin \omega t$$

$$|V_m| > V_0, \quad V_m < 0, \quad 0 < \omega t < 2\pi$$

$$f(y) = \frac{1}{\pi} \frac{1}{V_0 \sqrt{1 - \left( \frac{y - |V_m|}{V_0} \right)^2}}$$

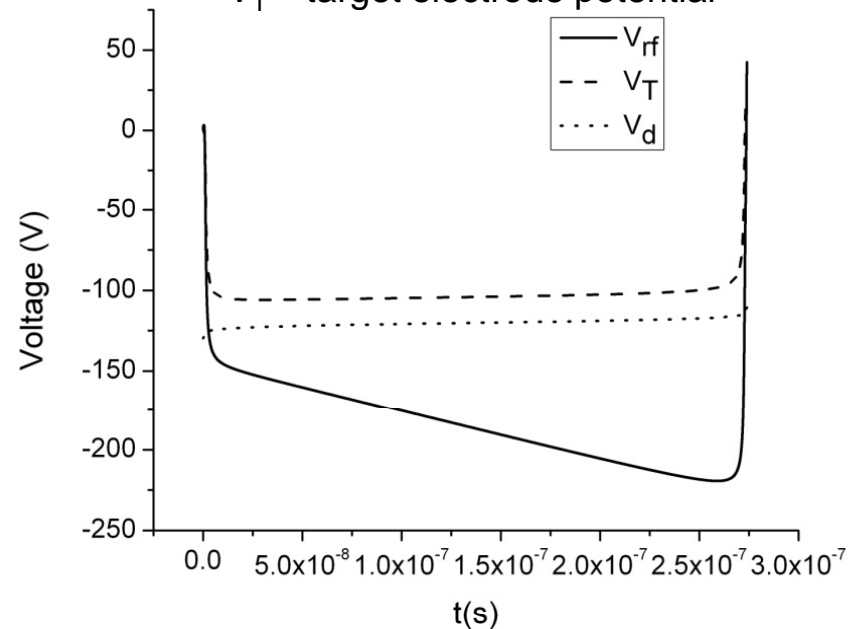
## Voltage Waveform Required to Obtain Desired IED (Inverse problem)

$C_b = 500$  pF,  $n_0 = 2 \times 10^{10}$  cm<sup>-3</sup>,  $T_e = 3$  eV,  
 $M = 40$  amu (Ar<sup>+</sup>),  $A_G/A_T = 20$ ,  $\omega\tau_i = 1$



*Desired* IED is Gaussian with specified peak and standard deviation.

$V_{rf}$  = rf voltage before blocking cap  
 $V_d$  = “damped” sheath potential  
 $V_T$  = target electrode potential

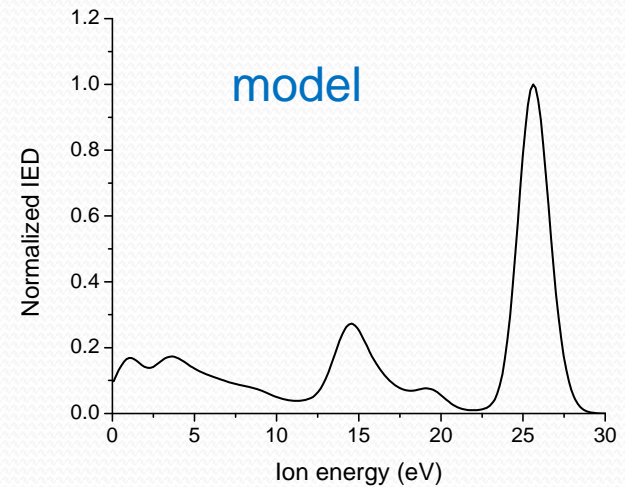
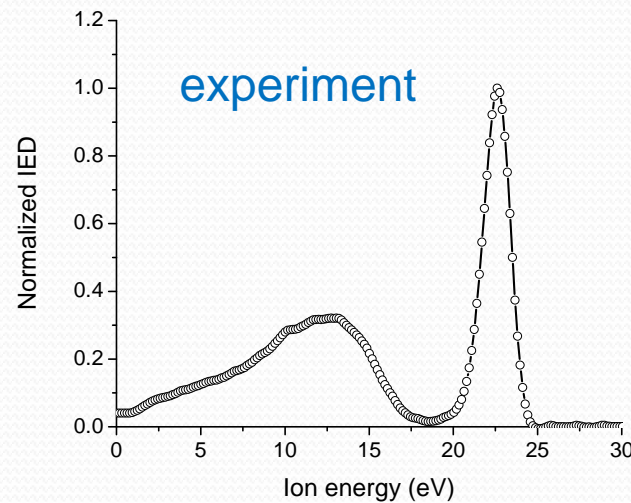


Required voltage waveform  $V_{rf}$  is a “rectangular” pulse with a slope. Slope is needed because of blocking capacitor charging.

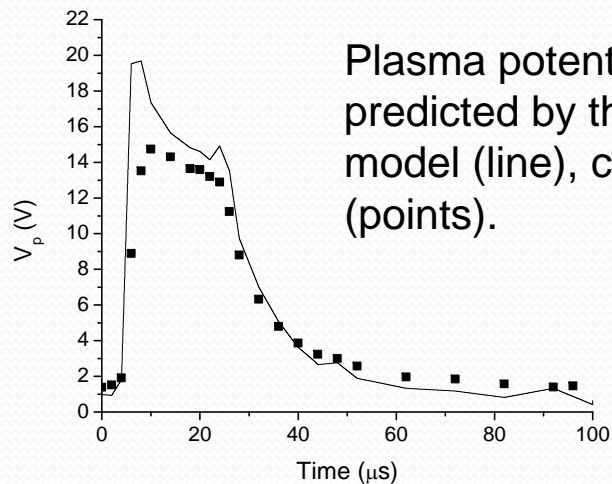
# Comparison of Semi-analytic Model with Experimental Data Argon Plasma

PS2-TuA9

PS+SS-WeM6



IEDs predicted by the semi-analytical model (right) compared to data<sup>#</sup> (left).



**Pulsed plasma, 10 KHz modulation, 20% duty, 14 mTorr, 120 W average power, 24.4 V DC bias applied in afterglow during  $\Delta t_b = 45-95 \mu s$ .**

<sup>#</sup> H. Shin et al., *PSST*, **20** 055001 (2011).



## Summary

- The energy distribution of ions bombarding the wafer can be tailored by applying voltage waveforms with special shapes (square wave, staircase, etc.).
- Both the location of the peak(s) of the IED and the fraction of ions under each peak can be controlled.
- Semi-analytic model can identify voltage waveforms that result in desired IED (inverse problem).
- Semi-analytic model executes much faster, but PIC simulation provides detailed information about both the IED and the EEDF upon application of bias.
- Experimental measurements were in good agreement with model and simulation predictions.



## Acknowledgements

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